Raffello Sanzio Scuola di Atene

High-energy physics with particle colliders: the present and the future

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Jide 7

- Physics results from LHC
 - Discovery of the 125 GeV Higgs boson
 - Measurement of mass and couplings; studies of spin & CP properties of this new particle
 - QCD, Electroweak, top, heavy-flavour
- Standard Model (SM): a complete theory
- New physics searches
 - Excess seen in diphoton production at $m_{\gamma\gamma} \sim 750 \text{ GeV}$
- Open questions in the Standard Model
- What Next?:
 - the LHC luminosity upgrade
 - Future accelerator facilities

Where are we today?

- *July 2012*: ATLAS and CMS announce the observation of a new, neutral particle with mass near 125 GeV, compatible with the Higgs boson predicted by the Standard Model
- The discovery is made searching for the Standard Model *bosonic* decays:
 - $-H \rightarrow \gamma \gamma$
 - $-H \rightarrow ZZ^* \rightarrow 41 (l=e,\mu)$
 - $-H \rightarrow WW^* \rightarrow lv lv (l=e,\mu)$
 - Recently, preliminary results from the $\sqrt{s=7,8}$ TeV run (*Run1*) based on the ATLAS and CMS combination have provided the observation of the H $\rightarrow \tau\tau$ decay mode as well. First observation of a *fermionic* Higgs boson decay: see:

https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2015-044/



Run Number: 189280, Event Number: 143576946 Date: 2011-09-14, 11:37:11 CET

EtCut>0.3 GeV PtCut>3.0 GeV Vertex Cuts: Z direction <1cm Rphi <1cm

Muon: blue Cells:Tiles, EMC $H \rightarrow ZZ \rightarrow 4I$

4







The distribution of the four-lepton invariant mass, $m_{4\ell}$. The signal expectation shown is for a mass hypothesis of $m_H=125$ GeV



CMS Experiment at the LHC, CERN Data recorded: 2012-May-13 20:08:14.621490 GMT Run/Event: 194108 / 564224000





1σ

2σ

Зσ

4σ

5σ

150



Background + Signal fit Sum of the m_{yy} distribution from the 7 and 8 TeV datasets, together with the data binned.

Best-fit signal strength, μ^{2} Signal significance = 5.7σ (5.7) expected)

atlas-conf-2015-044 slide 8 الملكة 125 GeV Higgs boson signal strength



125 GeV Higgs boson couplings



Spin/CP properties

- Study Higgs boson decay kinematic to extract spin/CP properties
- Spin 1 and 2 hypotheses (in many variants) are excluded at >95% C.L.



Higgs boson mass

- Combined mass through simultaneous fit to $H \rightarrow 41$ and $H \rightarrow \gamma \gamma$ datasets
- Scale accuracy at the level of 0.1%! ← spectrometer & calorimeter calibration important H→41 H→γγ combination



Uncertainty dominated by statistics, it will improve with more data from Run2

Higgs boson mass

- Combined mass through simultaneous fit to $H \rightarrow 41$ and $H \rightarrow \gamma \gamma$ datasets
 - Scale accuracy at the level of 0.1%! spectrometer & calorimeter
 - *m_H* was the only fundamental parameter missing to completely set the SM structure
 - We just entered a new era: precision tests are now needed in order to "discover" possible deviations from SM

	ATLAS	CMS
H → 41	$124.51 \pm 0.52 \pm 0.06$	$125.6 \pm 0.4 \pm 0.2$
Н→үү	$125.98 \pm 0.42 \pm 0.28$	$124.70 \pm 0.31 \pm 0.15$
Combination	$125.36 \pm 0.37 \pm 0.18$	$125.02 \ {}^{+0.26}_{-0.27} \ {}^{+0.14}_{-0.15}$
ATLAS + CMS	125.09 ± 0.21 (stat) ± 0.11 (scal	e) ± 0.02 (other) ± 0.01 (theory)

Uncertainty dominated by statistics, it will improve with more data from Run2



ATLAS SUSY Searches* - 95% CL Lower Limits

Status: March 2016

	Model	e, μ, au, γ	Jets	$E_{ m T}^{ m miss}$	∫ <i>L dt</i> [fb	⁻¹] Mass limit	$\sqrt{s} = 7, 8 \text{ TeV}$ $\sqrt{s} = 13 \text{ TeV}$	Reference
Inclusive Searches	$ \begin{array}{l} MSUGRA/CMSSM \\ \tilde{q}\tilde{q}, \tilde{q} \rightarrow q \tilde{\chi}_{1}^{0} \\ \tilde{q}\tilde{q}, \tilde{q} \rightarrow q \tilde{\chi}_{1}^{0} \\ (compressed) \\ \tilde{q}\tilde{q}, \tilde{q} \rightarrow q (\ell \ell \lfloor \ell \nu / \nu \nu) \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q \tilde{q} \tilde{\chi}_{1}^{+} \rightarrow q q W^{\pm} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q \tilde{\chi}_{1}^{+} \rightarrow q W^{\pm} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q (\ell \ell / \ell \nu / \nu \nu) \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q W Z \tilde{\chi}_{1}^{0} \\ GMSB (\ell NLSP) \\ GGM (hion NLSP) \\ GGM (higgsino-bino NLSP) \\ GGM (higgsino hion NLSP) \\ GGM (higgsino NLSP) \\ GGM (higgsino NLSP) \\ Gravitino LSP \end{array} $	$\begin{array}{c} 03 \ e, \mu/12 \ \tau \\ 0 \\ \text{mono-jet} \\ 2 \ e, \mu \ (\text{off-}Z) \\ 0 \\ 1 \ e, \mu \\ 2 \ e, \mu \\ 0 \\ 1\text{-}2 \ \tau + 0\text{-}1 \ \ell \\ 2 \ \gamma \\ \gamma \\ 2 \ e, \mu \ (Z) \\ 0 \end{array}$	2-10 jets/3 2-6 jets 1-3 jets 2-6 jets 2-6 jets 0-3 jets 7-10 jets 0-2 jets 2 jets 2 jets 2 jets mono-jet	b Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 3.2 20.3 3.2 3.3 20 3.2 20.3 20.3 2	\$\vec{q}\$.\$\vec{g}\$ 980 GeV \$\vec{q}\$ 610 GeV \$\vec{q}\$ 610 GeV \$\vec{q}\$ 820 GeV \$\vec{g}\$ 800 GeV \$\vec{g}\$ 900 GeV \$\vec{g}\$ 900 GeV \$\vec{g}\$ 900 GeV	$\begin{array}{cccc} \textbf{1.85 TeV} & \textbf{m}(\tilde{q}) = \textbf{m}(\tilde{g}) \\ & \textbf{m}(\tilde{k}^{0}) = 0 \text{ GeV}, \textbf{m}(1^{st} \text{ gen}, \tilde{q}) = \textbf{m}(2^{nd} \text{ gen}, \tilde{q}) \\ & \textbf{m}(\tilde{k}^{0}) = 0 \text{ GeV}, \textbf{m}(1^{st} \text{ gen}, \tilde{q}) = \textbf{m}(2^{nd} \text{ gen}, \tilde{q}) \\ & \textbf{m}(\tilde{q}) = 0 \text{ GeV} \\ & \textbf{m}(\tilde{k}^{0}) = 0 \text{ GeV} \\ \textbf{1.52 TeV} & \textbf{m}(\tilde{k}^{0}) = 0 \text{ GeV} \\ \textbf{1.6 TeV} & \textbf{m}(\tilde{k}^{0}) = 350 \text{ GeV}, \textbf{m}(\tilde{k}^{\pm}) = 0.5(\textbf{m}(\tilde{k}^{0}) + \textbf{m}(\tilde{g})) \\ \textbf{1.38 TeV} & \textbf{m}(\tilde{k}^{0}) = 0 \text{ GeV} \\ \textbf{1.6 TeV} & \textbf{m}(\tilde{k}^{0}) = 100 \text{ GeV} \\ \textbf{1.63 TeV} & tan\beta > 20 \\ \textbf{1.34 TeV} & cr(\textbf{NLSP}) < 0.1 \text{ mm} \\ \textbf{1.37 TeV} & \textbf{m}(\tilde{k}^{0}_{1}) = 950 \text{ GeV}, cr(\textbf{NLSP}) < 0.1 \text{ mm}, \mu < 0 \\ \textbf{m}(\tilde{k}^{0}_{1}) = 950 \text{ GeV}, cr(\textbf{NLSP}) < 0.1 \text{ mm}, \mu > 0 \\ \textbf{m}(\tilde{k}^{0}_{1}) = 430 \text{ GeV} \\ \textbf{m}(\tilde{d}) > 1.8 \times 10^{-4} \text{ eV}, \textbf{m}(\tilde{g}) = \textbf{n}(\tilde{g}) = 1.5 \text{ TeV} \end{array}$	1507.05525 ATLAS-CONF-2015-062 <i>To appear</i> 1503.03290 ATLAS-CONF-2015-062 ATLAS-CONF-2015-076 1501.03555 1602.06194 1407.0603 1507.05493 1507.05493 1507.05493 1503.03290 1502.01518
3 rd gen. <i>§</i> med.	$\begin{array}{l} \tilde{g}\tilde{g},\tilde{g} \rightarrow b \tilde{b} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g},\tilde{g} \rightarrow t \tilde{t} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g},\tilde{g} \rightarrow b \tilde{t} \tilde{\chi}_{1}^{+} \end{array}$	0 0-1 <i>e</i> ,μ 0-1 <i>e</i> ,μ	3 b 3 b 3 b	Yes Yes Yes	3.3 3.3 20.1	ē ē ē	1.78 TeV m(k ⁰ ₁)<800 GeV 1.76 TeV m(k ⁰ ₁)=0 GeV 1.37 TeV m(k ⁰ ₁)<300 GeV	ATLAS-CONF-2015-067 <i>To appear</i> 1407.0600
3 rd gen. squarks direct production	$ \begin{split} \tilde{b}_{1}\tilde{b}_{1}, \ \tilde{b}_{1} \to b\tilde{\chi}_{1}^{0} \\ \tilde{b}_{1}\tilde{b}_{1}, \ \tilde{b}_{1} \to \lambda\tilde{\chi}_{1}^{\pm} \\ \tilde{t}_{1}\tilde{t}_{1}, \ \tilde{t}_{1} \to b\tilde{\chi}_{1}^{\pm} \\ \tilde{t}_{1}\tilde{t}_{1}, \ \tilde{t}_{1} \to Wb\tilde{\chi}_{1}^{0} \text{ or } t\tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{1}, \ \tilde{t}_{1} \to Wb\tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{1}, \ \tilde{t}_{1} \to \tilde{t}_{1} \\ \tilde{t}_{1}\tilde{t}_{1} (natural GMSB) \\ \tilde{t}_{2}\tilde{t}_{2}, \ \tilde{t}_{2} \to \tilde{t}_{1} + Z \\ \tilde{t}_{2}\tilde{t}_{2}, \ \tilde{t}_{2} \to \tilde{t}_{1} + h \end{split} $	0 2 e, µ (SS) 1-2 e, µ 0-2 e, µ (0 m 2 e, µ (Z) 3 e, µ (Z) 1 e, µ	2 b 0-3 b 1-2 b 0-2 jets/1-2 nono-jet/c-ta 1 b 1 b 6 jets + 2 b	Yes Yes Yes Ves ag Yes Yes Yes Yes	3.2 3.2 1.7/20.3 20.3 20.3 20.3 20.3 20.3 20.3	b1 840 GeV b1 325-540 GeV c117-170 GeV 200-500 GeV c1 90-198 GeV 200-5715 GeV c1 90-245 GeV 745-785 GeV c1 90-245 GeV 200-610 GeV c1 290-610 GeV 200-610 GeV c2 320-620 GeV 320-620 GeV	$\begin{array}{c} m(\tilde{\chi}_{1}^{0}) \! < \! 100 \text{GeV} \\ m(\tilde{\chi}_{1}^{0}) \! = \! 50 \text{GeV}, m(\tilde{\chi}_{1}^{1}) \! = \! m(\tilde{\chi}_{1}^{0}) \! + \! 100 \text{GeV} \\ m(\tilde{\chi}_{1}^{0}) \! = \! 2\mathfrak{m}(\tilde{\chi}_{1}^{0}), m(\tilde{\chi}_{1}^{0}) \! = \! 55 \text{GeV} \\ m(\tilde{\chi}_{1}^{0}) \! = \! 1 \text{GeV} & 1506 \\ m(\tilde{\chi}_{1}^{0}) \! - \! 1 \text{GeV} & 1506 \\ m(\tilde{\chi}_{1}^{0}) \! < \! 85 \text{GeV} \\ m(\tilde{\chi}_{1}^{0}) \! > \! 150 \text{GeV} \\ m(\tilde{\chi}_{1}^{0}) \! = \! 0 \text{GeV} \\ m(\tilde{\chi}_{1}^{0}) \! = \! 0 \text{GeV} \end{array}$	ATLAS-CONF-2015-066 1602.09058 1209.2102, 1407.0583 08616, ATLAS-CONF-2016-00 1407.0608 1403.5222 1403.5222 1506.08616
EW direct	$ \begin{split} \tilde{\ell}_{L,R} \tilde{\ell}_{L,R}, \tilde{\ell} \to \ell \tilde{\chi}_1^0 \\ \tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \to \tilde{\ell} \nu(\ell \tilde{\nu}) \\ \tilde{\chi}_1^+ \tilde{\chi}_1^-, \tilde{\chi}_1^+ \to \tilde{\nu} \nu(\tau \tilde{\nu}) \\ \tilde{\chi}_1^+ \tilde{\chi}_2^0 \to \tilde{\ell}_1 \nu \tilde{\ell}_1 \ell(\tilde{\nu}\nu), \ell \tilde{\nu} \tilde{\ell}_L \ell(\tilde{\nu}\nu) \\ \tilde{\chi}_1^+ \tilde{\chi}_2^0 \to W \tilde{\chi}_1^0 Z \tilde{\chi}_1^0 \\ \tilde{\chi}_1^+ \tilde{\chi}_2^0 \to W \tilde{\chi}_1^0 h \tilde{\chi}_1^-, h \to b \tilde{b} / W W / \tau \\ \tilde{\chi}_2^+ \tilde{\chi}_3^0 \to \tilde{\chi}_3^0 \to \tilde{\ell}_R \ell \\ \end{split} $	$\begin{array}{c} 2 \ e, \mu \\ 2 \ e, \mu \\ 2 \ \tau \\ 3 \ e, \mu \\ 2 \ 3 \ e, \mu \\ 2 \ 3 \ e, \mu \\ 2 \ 3 \ e, \mu \\ \gamma \gamma \gamma e, \mu, \gamma \\ 4 \ e, \mu \\ . 1 \ e, \mu + \gamma \end{array}$	0 0 	Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3	$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	$\begin{split} & m(\bar{\chi}_{1}^{0}) \!=\! 0 \text{GeV} \\ & m(\bar{\chi}_{1}^{0}) \!=\! 0 \text{GeV}, m(\bar{\ell}, \bar{\nu}) \!=\! 0.5(m(\bar{\chi}_{1}^{+}) \!+\! m(\bar{\chi}_{1}^{0})) \\ & m(\bar{\chi}_{1}^{0}) \!=\! 0 \text{GeV}, m(\bar{\tau}, \bar{\nu}) \!=\! 0.5(m(\bar{\chi}_{1}^{+}) \!+\! m(\bar{\chi}_{1}^{0})) \\ & m(\bar{\chi}_{1}^{+}) \!=\! m(\bar{\chi}_{2}^{0}), m(\bar{\chi}_{1}^{0}) \!=\! 0, m(\bar{\ell}, \bar{\nu}) \!=\! 0.5(m(\bar{\chi}_{1}^{+}) \!+\! m(\bar{\chi}_{1}^{0})) \\ & m(\bar{\chi}_{1}^{+}) \!=\! m(\bar{\chi}_{2}^{0}), m(\bar{\chi}_{1}^{0}) \!=\! 0, sleptons decoupled \\ & m(\bar{\chi}_{1}^{+}) \!=\! m(\bar{\chi}_{2}^{0}), m(\bar{\chi}_{1}^{0}) \!=\! 0, sleptons decoupled \\ & m(\bar{\chi}_{2}^{0}) \!=\!\! m(\bar{\chi}_{2}^{0}), m(\bar{\chi}_{1}^{0}) \!=\! 0, m(\bar{\ell}, \bar{\nu}) \!=\!\! 0.5(m(\bar{\chi}_{2}^{0}) \!+\!\! m(\bar{\chi}_{1}^{0})) \\ & c\tau \!<\! 1 m \end{split}$	1403.5294 1403.5294 1407.0350 1402.7029 1403.5294, 1402.7029 1501.07110 1405.5086 1507.05493
Long-lived particles	Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\lambda}$ Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\lambda}$ Stable, stopped \tilde{g} R-hadron Metastable \tilde{g} R-hadron GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tau (\tilde{e}, \tilde{\mu}) + \tau$ GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_1^0$ $\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow eev/e\mu\nu/\mu\mu\nu$ GGM $\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow Z\tilde{G}$		1 jet - 1-5 jets - - - μ - ts -	Yes Yes - - Yes - -	20.3 18.4 27.9 3.2 19.1 20.3 20.3 20.3	$\begin{array}{c c c c c c c c c c c c c c c c c c c $	$\begin{array}{c} \mathfrak{m}(\tilde{k}_{1}^{*}){-}\mathfrak{m}(\tilde{k}_{1}^{0}){-}160~\text{MeV},~\tau(\tilde{k}_{1}^{*}){=}0.2~\text{ns}\\ \mathfrak{m}(\tilde{k}_{1}^{*}){-}\mathfrak{m}(\tilde{k}_{1}^{0}){-}160~\text{MeV},~\tau(\tilde{k}_{1}^{*}){<}15~\text{ns}\\ \mathfrak{m}(\tilde{k}_{1}^{0}){=}100~\text{GeV},~10~\mu\text{s}{<}\tau(\tilde{g}){<}1000~\text{s}\\ \mathfrak{m}(\tilde{k}_{1}^{0}){=}100~\text{GeV},~\tau{>}10~\text{ns}\\ 10{<}\tan\theta{<}50\\ 1{<}\tau(\tilde{k}_{1}^{0}){<}3~\text{ns},~\text{SPS8}~\text{model}\\ 7{<}c\tau(\tilde{k}_{1}^{0}){<}740~\text{mm},~\mathfrak{m}(\tilde{g}){=}1.3~\text{TeV}\\ 6{<}c\tau(\tilde{k}_{1}^{0}){<}480~\text{mm},~\mathfrak{m}(\tilde{g}){=}1.1~\text{TeV}\\ \end{array}$	1310.3675 1506.05332 1310.6584 <i>To appear</i> 1411.6795 1409.5542 1504.05162 1504.05162
RPV	$ \begin{array}{l} LFV \ pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e\mu/e\tau/\mu_{1} \\ Bilinear \ RPV \ CMSSM \\ \tilde{\lambda}_{1}^{+} \tilde{\lambda}_{1}^{-}, \tilde{\lambda}_{1}^{+} \rightarrow W \tilde{\lambda}_{1}^{0}, \tilde{\lambda}_{1}^{0} \rightarrow ee\tilde{v}_{\mu}, e\mu\tilde{v} \\ \tilde{\lambda}_{1}^{+} \tilde{\lambda}_{1}^{-}, \tilde{\lambda}_{1}^{+} \rightarrow W \tilde{\lambda}_{1}^{0}, \tilde{\lambda}_{1}^{0} \rightarrow \tau \tau \tilde{v}_{e}, e\tau\tilde{v} \\ \tilde{s}\tilde{g}, \tilde{g} \rightarrow qqq \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_{1}^{0}, \tilde{\lambda}_{1}^{0} \rightarrow q\bar{q}q \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{t}_{1}\tilde{t}, \tilde{t}_{1} \rightarrow bs \\ \tilde{t}_{1}\tilde{t}_{1}, \tilde{t}_{1} \rightarrow bl \end{array} $	$\begin{array}{ccc} & e\mu, e\tau, \mu\tau \\ & 2 \ e, \mu \ (\text{SS}) \\ \phi_e & 4 \ e, \mu \\ \phi_\tau & 3 \ e, \mu + \tau \\ & 0 \\ & 2 \ e, \mu \ (\text{SS}) \\ & 0 \\ & 2 \ e, \mu \end{array}$	- 0-3 b - 6-7 jets 6-7 jets 0-3 b 2 jets + 2 b 2 b	- Yes Yes - - Yes - Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3	$ \begin{array}{c} \tilde{v}_{\tau} \\ \tilde{q}, \tilde{g} \\ \tilde{\chi}_{1}^{*} \\ & 760 \text{ GeV} \\ \tilde{\chi}_{1}^{*} \\ & 450 \text{ GeV} \\ \tilde{g} \\ \tilde{g} \\ & 917 \text{ GeV} \\ \tilde{g} \\ & 980 \text{ GeV} \\ \tilde{g} \\ & 880 \text{ GeV} \\ \tilde{t}_{1} \\ & 320 \text{ GeV} \\ \hline{t}_{1} \\ & 0.4-1.0 \text{ TeV} \\ \end{array} $	1.7 TeV $\lambda'_{311}=0.11, \lambda_{132/133/233}=0.07$ 1.45 TeV $m(\tilde{q})=m(\tilde{g}), c\tau_{LSP}<1 \text{ mm}$ $m(\tilde{k}_{1}^{0})>0.2\times m(\tilde{k}_{1}^{+}), \lambda_{121}\neq 0$ $m(\tilde{k}_{1}^{0})>0.2\times m(\tilde{k}_{1}^{+}), \lambda_{133}\neq 0$ BR(t)=BR(b)=BR(c)=0% $m(\tilde{k}_{1}^{0})=600 \text{ GeV}$ $BR(\tilde{t}_{1}\rightarrow be/\mu)>20\%$	1503.04430 1404.2500 1405.5086 1405.5086 1502.05686 1502.05686 1404.2500 1601.07453 ATLAS-CONF-2015-015
Other	Scalar charm, $\tilde{c} \rightarrow c \tilde{\chi}_1^0$	0	2 c	Yes	20.3	č 510 GeV	m(₹10)<200 GeV	1501.01325
*Onl	y a selection of the availa	ble mass limi	ts on new	/	1	D ⁻¹	1 Mass scale [TeV]	,

*Only a selection of the available mass limits on new states or phenomena is shown.

Mass scale [TeV]

ATLAS Preliminary $\sqrt{s} = 7, 8, 13 \text{ TeV}$

Searches for yy resonances

• Select events with two high- p_T isolated photons – plot the $\gamma\gamma$ invariant mass, $m_{\gamma\gamma}$



Searches for high-mass yy resonances

ATLAS

- Largest deviation from B-only hypothesis m_X ~750 GeV, Γ_X ~45GeV(6%)
- Local significance=3.9σ
- Global significance=2.0σ



- Largest excess observed for $m_X = 760$ GeV and $\Gamma/m = 1.4 \times 10^{-2}$
- Similar significance for narrow-width hypothesis.
- "Global" significance < 1σ.

What Next?

What Next?

- Many open points in high-energy physics, still. Among these:
 - Is the Standard Model an *effective theory* of more general fundamental theory, or is it rather a *sector of the final fundamental theory*?
 - Is the 125 GeV Higgs boson the fundamental scalar predicted by SM, or is it explained by an extend theory?
 - What's the nature and the origin of **Dark Matter**?
 - *Hierarchy*, problem (aka *fine-tuning* or *naturalness*), calling for new particles to stabilize quantum corrections to the Higgs boson mass (but it is a real problem??)
 - *Flavour problem*: what's the explanation of the large matter over antimatter dominance in the universe?
 - Neutrino masses
 - (many others, not listed here)

What's Next – an example

arXiv:1206.3560v3

Rick S. Gupta^{a,b,c}, Heidi Rzehak^{a*}, James D. Wells^{a,b}

	ΔhVV	$\Delta h ar{t} t$	$\Delta h \overline{b} b$
Mixed-in Singlet	6%	6%	6%
Composite Higgs	8%	tens of $\%$	tens of $\%$
Minimal Supersymmetry	< 1%	3%	$10\%^a, 100\%^b$
LHC Run1 experiment	~15%	~40%	~30%
LHC 14 TeV 300 fb ⁻¹ (1 exp.)	~6%	~15%	~12%

- 14 TeV & 300 fb⁻¹ projections based on what we learnt from early 2012 analyses
 - Since then, more channels have been included in the 125 GeV Higgs boson data analyses

Which tools?

- LHC real: present and near future
- The LHC Luminosity Upgrade, HL-LHC the concrete *mid-term future*
- ... and then the *far future* open options:
- The LHC energy upgrade, **HE-LHC**
- Linear e⁺e⁻ colliders
- Circular e⁺e⁻ colliders
- Photon colliders
- Very high-energy hadron colliders
- Muon collider(s)

Next steps -1-

- LHC is the only high-energy particle collider in operation today at it will be the only one for many years to come!
- All future projects in HEP should be based on the outcome from LHC (√s≅14 TeV; L≅300 fb⁻¹)
- We have two possible situations after LHC (300 fb⁻¹)
 - 1. LHC will not provide direct observation of New Physics at O(1 TeV) scale, or
 - 2. New Physics will be observed at LHC

Next steps -2-

- In **both cases** the luminosity upgrade of LHC, <u>**HL-**</u> <u>**LHC**</u>, is the most realistic and sensible project to perform in the short term:
- 1. We do not find New Physics at the LHC →
 - a. Precision studies of the 125 GeV Higgs boson sector
 - b. Precision Measurements of Standard Model in general
 - **c.** Continue the search for New Physics, in particular for rare processes. Deviations from SM will indicate presence of New Physics not too far away from the energy scale of the LHC
 - All of this represents <u>essential input</u> for future collider experiments!

Next steps -3-

• Need to change mentality in HEP: it is not true that

New Physics =(only) discovery of new particles!

- → New Physics is in the observation of New Phenomena
- LHC has delivered the (first?) new phenomenon: a neutral scalar of 125 GeV mass
 - If elementary, it would be the first point-like scalar ever seen in nature!
 - Must look at it closely, it represents a portal to a completely unknown new territory

Next steps -4-

2. We find new physics at the LHC \rightarrow it will redesign the priorities of HEP.

Need to understand what this new phenomena would be, and with more data from LHC/HL-LHC would be important to a) to validate with statistics the nature of the effects seen, and b) to study the associated properties

In addition perform the Precision Measurements already mentioned

- Example (from these days...): the 750 GeV γγ excess, if confirmed to be of physics nature by the data we're collecting this year
- All of this represents <u>essential input</u> for future collider experiments

The High-Luminosity LHC



The High Luminosity LHC

LHC / HL-LHC Plan





The High Luminosity LHC

LHC / HL-LHC Plan





the CERN Council adopted the European Strategy for Particle Physics in Brussels on 30 May 2013, its first priority was agreed to be "[...] Europe's top priority should be the exploitation of the full potential of the LHC, including the high-luminosity upgrade of the machine and detectors with a view to collecting ten times more data than in the initial design, by around 2030 [...]".

The High Luminosity LHC

LHC / HL-LHC Plan





- Full exploitation of LHC is top priority in Europe & US for high energy physics
- Operate HL-LHC with 5 (nominal) to 7.5 (ultimate) $\times 10^{34}$ cm⁻²s⁻¹ to collect 3000 fb⁻¹ in ~ ten years.

The challenge of HL-LHC

- The challenge of HL-LHC, both for detectors and physics, is the **event pile-up**
 - $n = (\sigma L)/(n_B f_r)$ with $\sigma \sim 81$ mb, L=7.5 10³⁴ cm⁻² s⁻¹, $n_B = 2808$ and $f_r = 11450$ Hz $\rightarrow n \sim 200$ events per bunch crossing



Detector Upgrade

- Some of present the sub-systems will either not survive the high luminosity environment (radiation damage) or not function efficiently because of the increased data rates (detector occupancy and event pile-up).
- → important to upgrade the present ATLAS and CMS detectors to cope with the HL-LHC environment!
 - Entirely **new inner tracking systems** to measure charged particles will be required, and
 - the energy-measuring **calorimeters** will also need partial replacement, in the endcap region for CMS and possibly in the more forward region for ATLAS

Physics at HL-LHC

Physics at HL-LHC

- Lots of physics simulation performed since 2012 (Euopean Strategy, Cracow), and revisited afterwards. These includes:
- 125 GeV Higgs boson rates & couplings

 ttH production, rare decays such as H→μμ, Zγ, J/ψγ, …
- 125 GeV Higgs boson 2HDM interpretation assuming SM production and decays
- 125 GeV Higgs width, CP property studies
- 125 GeV Higgs *Invisible* decays and portal to Dark Matter
- BSM Higgses decays to bosons
- Double Higgs production and self-coupling
- Vector Boson Scattering
- New Physics: SUSY process: electroweak production studies, third generation searches
- New Physics: mono-Higgs, dilepton resonances, di-top resonances, top quark FCNC

Higgs boson sector

- Many aspects must be studied in detail:
 - Production/decay rates and rare modes, in particular
 - the Yukawa coupling with the top quark, ttH
 - the $H \rightarrow \mu^+ \mu^-$ and other rare decays
 - the Higgs boson couplings in general
 - the natural width $\Gamma_{\rm H}$ (SM: ~ 4.2 MeV)
 - Differential distributions: $d\sigma/dp_T$, $d\sigma/d\eta,$...
 - Spin/CP properties
 - Higgs selfcoupling
 - Search for partners of the 125 GeV Higgs boson

$H \rightarrow ZZ^* \rightarrow 41$

The large event yield of high mass resolution processes $H \rightarrow \gamma \gamma$ and $H \rightarrow ZZ^* \rightarrow 41$ will provide large statistics and high-purity samples for precision property and differential distribution measurements



Four lepton mass distributions for the full Phase-II luminosity of 3000 fb⁻¹ for the signal, $H \rightarrow ZZ \rightarrow 41$, and for the irreducible $ZZ \rightarrow 41$ background.





Rare decays

H→μμ

- Allows direct study of coupling to two different leptons
- Signal significance $> 7\sigma$

ATL-PHYS-PUB-2013-014

\mathcal{L} [fb ⁻¹]	300	3000
Signal significance	2.3σ	7.0σ
$\Delta \mu / \mu$	46%	21%

• CMS: revised projection, expect 5% on coupling measurement at HL-LHC

ATL-PHYS-PUB-2014-006

Н→Ζγ

- In Standard Model, this decay proceeds entirely via loops predominantly involving heavy charged particles
- \rightarrow sensitive to possible new physics
- $\sim 4\sigma$ evidence of the SM decay is possible at HL-LHC



CMS H→µµ coupling precision improves from 8% to 5% with Phase II upgrade

Precision on signal strength

NS Drag			1 atra	noth	suite
PIEC	ISIOII OII S	agna	ISUE	ngu	_
channel	Prec. (%) 100 fb ⁻¹	Prec. (%)	300 fb ⁻¹	Prec. (%)	3000 fb ⁻¹
ttH H→γγ	~65	38	36	17	12
ttH H \rightarrow ZZ* \rightarrow 41	~85	49	48	20	16
VBF H → γγ	~80	47	43	22	15
VBF H \rightarrow ZZ* \rightarrow 41	~60	36	33	21	16
Н→μμ	~70	39	38	16	12
Η→ττ	~18	14	8	8	5
H→bb	~20	14	11	7	5
Η→γγ	~15	12	6	8	4
H → 41	~15	11	7	9	4
H → 41	~15	11	7	7	4
	t		My pers	onal estima	ites

ATLAS: experimental & theory uncertianties; only exp. uncertainty current exp.l & theory uncertianties; exp. uncertainty $\propto 1/\sqrt{L}$ and $\frac{1}{2}$ theory unc. CMS:

ATLAS assumed luminosity uncertainty: 3%

From signal rates to Higgs couplings

• The cross section times branching ratio for initial state *i* and final state *f* is given by

$$\sigma \cdot Br(i \to H \to f) = \frac{\sigma_i \cdot \Gamma_f}{\Gamma_H}$$

- The total width Γ_H is too narrow to measure at LHC
 - Assume it is the sum of the visible partial widths no additional invisible modes
- Cross sections and branching ratios scale with $\kappa^2~(\not\to\Delta\kappa\sim0.5~\Delta\mu)$
- Gluon and photon couplings can be assumed to depend on other SM couplings, or to be independent to allow for new particles in the loop



Higgs Couplings

- VH->bb included in ATLAS, updates for H->Zγ, VH/ttH->γγ (*)
- No BSM Higgs decay modes assumed
- –Comparable numbers for $\kappa_W, \kappa_{Z_1}, \kappa_{t_1}$ and κ_{γ} between the two experiments
- Couplings can be determined with 2-10% precision at 3000fb⁻¹ (for CMS Scenario 2)

		κ _γ	κ _w	κ _z	Kg	к _b	<mark>К</mark> t	κ _τ	κ _{Ζγ}	κ _μ
300fb ⁻¹	ATLAS	[9,9]	[9,9]	[8,8]	[11,14]	[22,23]	[20,22]	[13,14]	[24,24]	[21,21]
300fb ⁻¹	CMS	[5,7]	[4,6]	[4,6]	[6,8]	[10,13]	[14,15]	[6,8]	[41,41]	[23,23]
3000fb ⁻¹	ATLAS	[4,5]	[4,5]	[4,4]	[5,9]	[10,12]	[8,11]	[9,10]	[14,14]	[7,8]
3000fb ⁻¹	CMS	[2,5]	[2,5]	[2,4]	[3,5]	[4,7]	[7,10]	[2,5]	[10,12]	[8,8]

CMS: <u>http://arxiv.org/abs/1307.7135</u>

ATLAS: [no theory uncert., full theory uncert.]CMS: [Scenario 2, Scenario1]

(*) ATLAS documents: ATL-PHYS-PUB-2014-011 ATL-PHYS-PUB-2014-006 ATL-PHYS-PUB-2014-012 ATL-PHYS-PUB-2014-016

Coupling fit with L=300 and 3000 fb⁻¹

Measurement accuracy <u>per experiment</u>! (model dependent)

Coupling modifier	300 fb ⁻¹	3000 fb ⁻¹	
$k_{W,Z}, k_{\gamma,}$	6%	3%	
k _b	12%	5%	down-quark type
k _t	15%	7%	Top Yukawa coup.
$\mathbf{k}_{ au}$	10%	5%	lepton coupling
k _u	22%	7%	2 nd generation

based on the results presented in this talk

CMS: <u>http://arxiv.org/abs/1307.7135</u> ATLAS: ATL-PHYS-PUB-2014-016

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Model independent scenario:

based on coupling ratio measurements, in backup

Higgs self-coupling -1-

- Never seen self-interacting fundamental particles!
- Is the Higgs potential so "simple" as the one described by the Standard Model?
 - $-V(\phi) = \mu^{2}(\phi^{+}\phi) + \lambda(\phi^{+}\phi)^{2} \quad \lambda = m_{H}/\nu^{2}; \quad \mu^{2} = (1/2)m_{H}^{2}$
 - This is not the case in several BSM implementations → experimental verification is essential, e.g. measuring HH production
 - → Any deviation will reveal New Physics!
- The parameter λ can be probed studying the production of HH from the self-coupling process



Higgs self-coupling -1con't-

- Never seen self-interacting fundamental particles!
- Is the Higgs potential so "simple" as the one described by the Standard Model?

$$-V(\phi) = \mu^{2}(\phi^{+}\phi) + \lambda(\phi^{+}\phi)^{2} \quad \lambda = m_{H}/v^{2}; \quad \mu^{2} = (1/2)m^{2}_{H}$$

- This is not the case in several BSM implementations → experimental verification is essential, e.g. measuring HH production
- \rightarrow Any deviation will reveal New Physics!
- Example of New Physics processes:



 $hh \rightarrow bb\gamma\gamma$

Higgs self-coupling -2-

- For LHC, HL-LHC, several hh final states are already under study:
 - hh→ bbγγ
 - $-hh \rightarrow bb\tau\tau$
 - hh→ bbWW
 - $-hh \rightarrow \gamma \gamma WW$



process ATLAS	Expected events in 3000 fb ⁻¹
SM HH→bbγγ	8.4±0.1
bbyy	9.7 ± 1.5
ссүү, bbүj, bbjj, jjүү	24.1 ± 2.2
top background	3.4 ± 2.2
ttH(γγ)	6.1 ± 0.5
Z(bb)H(γγ)	2.7 ± 0.1
bbH(үү)	1.2 ± 0.1
Total background	47.1 ± 3.5
S/VB (barrel+endcap)	1.2

- Combination of HH \rightarrow bb $\tau\tau$ and HH \rightarrow bb $\gamma\gamma$ yields a significance of 1.9 standard deviations; measurement accuracy: 54% (per experiment)
- Measurement of HH production at the HL-LHC at the level of 30% should be possible



Expected 5 or Discoverv

• Naturalness motivates stop/sbottom searches where the third family squarks are lightest **CMS** Simulation √s = 14 TeV



5σ discovery, simplified model	300 fb ⁻¹	3000 fb ⁻¹
stop mass from direct production [ATLAS]	Up to 1.0 TeV	Up to 1.2 TeV
gluino mass with decay to stop [CMS]	Up to 1.9 TeV	Up to 2.2 TeV
sbottom mass from direct production [ATLAS]	Up to 1.1 TeV	Up to 1.3 TeV

The post-LHC era

• Many ideas on the table – two main groups:

-Higgs Factories:

- Electron-Positron colliders
 - Circular: FCC-ee, CepC
 - Linear: ILC, CLIC
- Photon collider: SAPPHiRE
- Muon collider

-Very High Energy hadron colliders

• Circular proton collider: HE-LHC, FCC-hh, CppC

Higgs boson production at lepton colliders



- In e+e- colliders, the dominant H production goes with the associated production channel e+e_ → ZH
 - It does require $E_{CM} > m_Z + m_H = 216$ GeV

- The s-channel allows direct lineshape scan, but the production rate is too
 small because the too small Yukawa coupling of H to electrons
- In µ+µ- colliders, the H production in the s-channel is significant
 - muon collider can produce in the schannel $(m_{\mu}/m_{e})^{2}$ =4.3×10⁴ more Higgs bosons
 - It does require $E_{CM} \sim m_H = 125 \text{ GeV}$
 - The only way for direct $\Gamma_{\rm H}$ measurement

FCC-ee/hh/eh

Description	value	Lake Geneva
Circumference [km]	(80) 100	
centre-of-mass energy e+e– [GeV]	91, 160, 240, 350	LEP/LHC Sure a state of the sta
Luminosity [10 ³⁴ cm ⁻² s ⁻¹]	5.1 @ 240 GeV	80-100 km tunnel
Synchrotron radiation loss [GeV/turn]	1.67 @ 240 (→ 50 MW/ beam)	LEGEND LHC tunnel HE_LHC 80km option potential shaft location

- Circular collider at CERN to collide
 - electron-positron
 - proton-proton (see later)
 - electron-proton (not discussed in this seminar)
 beams

CEPC

• electron-positron circular collider in China

http://cepc.ihep.ac.cn/intro.html updated on 3 August 2014

Description	value
Circumference [km]	53.6
Centre-of-mass energy [GeV]	240
Luminosity [10 ³⁴ cm ⁻² s ⁻¹]	1.80
Synchrotron radiation loss [GeV/turn]	3.01 (→ 50 MW/ beam)



Schematic of 50 or 100 km tunnel f or a Chinese electronpositron Collider (CEPC) and Super proton-proton Collider in Qinhuangdao (秦皇岛), with great geological conditions and strong support from the local municipal government, could be the potential site.



CDR in: https://edms.cern.ch/ui/file/1234244/7/CERN-2012-007.pdf

Description	500 GeV	3 TeV
Centre-of-mass energy [TeV]	0.5	3
Luminosity [10 ³⁴ cm ⁻² s ⁻¹]	2.3	5.9
Accelerator gradient [MV/ m]	80	100
Total Length [km]	13.0	48.4



Fig. 2.2: Map showing a potential location for the CLIC accelerator complex



ILC Facility Design (Technical Design Report 2014)



31.5 MV/m +/-20%

 $Q_0 = 1E10$



E gradient in SCRF

acc. cavity

Site Specific design: the "preferred" site⁵²



Muon collider



- The determination of the width $\Gamma_{\rm H}$ is be crucial in the determination of the nature of the particle and the underlying physics. The SM prediction is only ~ 4 on MeV, a formidable task!
- Assuming an instantaneous luminosity of 0.6×10³² cm⁻²s⁻¹, ~ 4400 (SM) Higgs bosons /year are produced (per experiment)



FCC-hh / HE-LHC

- FCC-hh: hadron collider in the FCC-ee/TLEP tunnel
 - $\sqrt{s} \sim 100$ TeV using 16T Nb₃Sn magnets
 - $\sqrt{s} \sim 50$ TeV using 8.4T NbTi magnets (LHC dipoles)

- HE-LHC: new hadron collider in the LHC tunnel
 - Needs new dipoles to increase the collision energy
 - For a 16T Nb₃Sn magnet we would get ~30 TeV



Example: Higgs boson couplings

Uncertainties	HL-LHC	μ Collider	CLIC 350 ~.5 ab ⁻¹	ILC 250 ~1.1 ab ⁻¹	CEPC 240 ~2 ab ⁻¹	FCC-ee 10 ab ⁻¹ 240 2.6 ab ⁻¹ 350	
m _H [MeV]	40?	0.06	40	30	5.5	8	
Γ _H [%]	[indirect]	4.0	5.0	5.4	4.4	1.0	
g _{HZZ} [%]	2	-	0.8	0.6	0.40	0.15	
g _{HWW} [%]	2	2.2	1.8	2.3	1.9	0.2	
g _{Hbb} [%]	4	2.3	2.0	2.5	2.1	0.4	
$g_{H\tau\tau}$ [%]	4	5	3.7	2.7	2.3	0.5	
g _{Hcc} [%]	-	-	3.2	3.2	2.7	0.7	
$\mathbf{g}_{\mathbf{H}\gamma\gamma}$ [%]	2	10	-	8.2	7.4	1.5	
g _{Hgg} [%]	3	-	3.6	3	2.4	0.8	
g _{Htt} [%]	5	-	-	-	-	-	
g _{Ημμ} [%]	5	2.1	-	-	14	6.2	
g _{HHH} [%]	30 →	-	-	-	-	-	
	ATLAS+CMS combination (my estimate)	Taken from M. Klute LCWS 2015	M. Pandurovic, Gomel School 2015	J. Tian, ICHEP2014	main_preCDR.	arXiv:1308.6176v3	

Conclusions

- With the discovery of the 125 GeV Higgs boson, a new era in high-energy physics is just opened!
- The Standard Model is a complete theory now:
 - test its validity at the highest possible theoretical and experimental precision
 - It does not explain all observations in nature: new physics must exist at some energy scale
 - The Higgs sector is experimentally completely new and its investigation may reveal precious information on where new physics could be
- Direct searches for new objects and new phenomena are essential
- LHC and its luminosity upgrade, HL-LHC, can shed light on the nature of the 125 Higgs boson at the level of O(few %) on couplings, and search for new particles at the O(few TeV) scale
- Decision on the next future collider(s) cannot be taken until we'll not first see what the outcome from LHC
 - Appointment with the European Strategy (symposium in 2018/19)



The LHC is almost back! ③

5:30 am on Friday 29 April Electrical perturbation

- Cause: short circuit caused by *fouine* on 66kV transformer in point 8
- Transformer connections damaged

Comments (05-May-2016 09:26:07)

66kV transformer @ P8 back in service preparing for test without beam till ~ 12:00

(morning meeting @ 9h during the long WE)

Comments (05–May–2016 14:08:30)

precycling

next: cycle with pilot beam

next injection in \sim 1 hour



backup

The ATLAS detector upgrade



ATLAS & CMS Upgrade

- Pileup mitigation is a critical element of the design
- More details on:
 - ATLAS [CERN-LHCC-2015-020]
 - & CMS [CERN-LHCC-2015-019]
 - "Scoping Documents" showing impact of different cost scenarios on physics performance, and on:
 - ATLAS Phase II LoI [CERN-LHCC-2012-022],
 - CMS Technical Proposal [CERN-LHCC-2015-010]

Example of pileup mitigation



Ratio of the number of reconstructed jets with the *PUPPI*¹ algorithm to the number of reconstructed jets matched to particle level jets for different p_T ft bins.

1. [http://arxiv.org/pdf/1407.6013v2.pdf]

The resolutions of the x and y components of E_T^{miss} in the three scoping scenarios for samples of tt events with μ =200. The resolutions are shown as a function of the scalar sum of total transverse energy in the event, with MC statistical uncertainties.



k-framework: *slide* 63 determination of coupling deviations

Coupling to fermions: $g_F = \kappa_F rac{\sqrt{2m_F}}{v}$ Coupling to bosons: $g_V = \kappa_V rac{2m_V^2}{v}$

Write rates in terms of k_i parameters The index *i* is related to each individual elementary particle

- Production: $\sigma_i \sim \kappa_i^2 \sigma_i^{SM}$
- Decay: $\Gamma_i \sim \kappa_i^2 \Gamma_i^{SM}$
- Total Width: $\Gamma_{\rm H} = \Sigma_{\rm i} \kappa_{\rm i}^2 \Gamma_{\rm i}^{\rm SM}$

Total width not measured at LHC (upper limits can be derived from off-shell production studies)

Assumptions:

- Only one Higgs boson
- Only scalar modifications of the coupling strength: kinematics as in SM
- $J^P = \theta^+$



Example:

Interference in $H \rightarrow \gamma \gamma$, $gg \rightarrow H$, ...: \rightarrow some relative sign-ambiguities solved

Higgs Couplings



- Most generic model: remove the assumption on the total width
 - Only ratios of the coupling scale factors can be determined at LHC
 - Use given process as a reference

ATL-PHYS-PUB-2015-024

- Total Higgs boson width can be probed at LHC studying the comparison of on-shell and off-shell Higgs boson production
- Recent study by ATLAS indicates that, under a number of assumptions, the (SM) Higgs boson width can be estimated with a systematic uncertainty of ~2 MeV:

 $\Gamma_{H}^{(L2)} = 4.2^{+1.5}_{-2.1}$ MeV (stat+sys).

Higgs Couplings





N³LO ggF production cross section



C. Anastasiou et al.: <u>http://arxiv.org/abs/1107.0683</u>, <u>http://arxiv.org/abs/1403.4616</u>, <u>http://arxiv.org/abs/1411.3584</u>, <u>http://arxiv.org/abs/1503.06056</u>

FIG. 2: Scale variation of the gluon fusion cross-section at all perturbative orders through $N^{3}LO$.

- For renormalisation and factorisation scales equal to half the Higgs mass, the N³LO corrections are of the order of +2.2%. The total scale variation at N³LO is 3%, reducing the uncertainty due to missing higher order QCD corrections by a factor of three.
- However the procedure to transfer the scale variations to theory uncertainty on the cross section is not firmly established (why considering a factor 2 and ½, for example?) [1] http://arxiv.org/abs/1409.5036 [2] http://arxiv.org/abs/1307.1843

67 Effects of theory uncertainties

- Theoretical uncertainties limit the achieved precision
- Reducing the theoretical uncertainties is a worthwhile endeavor

CMS:

Scenario 1 No theory uncertainty

CMS Projection



ATLAS: Deduced size of theory uncertainty to increase total uncertainty by <10% for 3000 fb⁻¹

Scenario	Status					
	2014	by $\leq 10\%$ for 3000 fb ⁻¹				
Theory uncertainty (%)	[10–12]	К _{gZ}	$\lambda_{\gamma Z}$	λ_{gZ}	$\lambda_{\tau Z}$	λ_{t_i}
$gg \rightarrow H$		1				
PDF	8	1.3	-	-	-	-
incl. QCD scale (MHOU)	7	1.1	-	-	-	-
p_T shape and $0j \rightarrow 1j$ mig.	10–20	-	1.5–3	-	-	-
$1j \rightarrow 2j$ mig.	13–28	-	3.3–7	-	-	-
$1j \rightarrow VBF 2j mig.$	18–58	-	-	6–19	-	-
VBF $2j \rightarrow VBF 3j$ mig.	12–38	-	-	-	6–19	-
qqH						
PDF	3.3	-	-	2.8	-	-
tīH						
PDF	9	-	-	-	-	3
incl. QCD scale (MHOU)	8	-	-	-	-	2

See also talk from W. Verkerke, this conference

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Interpretation for BSM Higgs sector

- Use expected accuracy in the couplings k_i of the 125 GeV Higgs boson
- Same decays as in the SM are assumed (b-associated production included as a correction)



Direct BSM Higgs searches: example $H \rightarrow ZZ \rightarrow 41$

- The 4l final state has small cross section but is clean and well reconstructed
- A heavy SM-like Higgs boson decaying to 4l occurs in several extensions of the scalar sector (2HDM,EWK singlet)
- Limits improve by a factor \sim 3 with 3000 fb⁻¹ wrt 300 fb⁻¹
- Similar sensitivity for ATLAS and CMS (~0.01-0.1 fb)



CLIC

Table 9.1: Key parameters of the 500 GeV and 3 TeV designs.

Description [units]	Parameter	500 GeV	3 TeV
Centre-of-mass energy [TeV]	$E_{ m cms}$	0.5	3.0
Repetition frequency [Hz]	$f_{\rm rep}$	50	50
Number of bunches per train	$n_{\rm b}$	354	312
Bunch separation [ns]	Δt	0.5	0.5
Accelerating gradient [MV/m]	G	80	100
Total luminosity [10 ³⁴ cm ⁻² s ⁻¹]	$L_{\rm total}$	2.3	5.9
Luminosity within 1% of $E_{\rm cms}$ [10 ³⁴ cm ⁻² s ⁻¹]	$L_{0.01}$	1.4	2.0
No. of photons per electron	n_{γ}	1.3	2.1
Average energy loss due to beamstrahlung	$\Delta E/E$	0.07	0.28
No. of coherent pairs per bunch crossing	$N_{\rm coh}$	200	6.8×10^{8}
Energy of coherent pairs per bunch crossing [TeV]	$E_{ m coh}$	15	2.1×10^{8}
No. of incoherent pairs per bunch crossing	nincoh	8×10^4	0.3×10^{6}
Energy of incoherent pairs per bunch crossing [GeV]	$E_{ m incoh}$	3.6×10^{6}	2.3×10^{7}
Hadronic events per crossing	<i>n</i> _{had}	0.3	3.2

• p